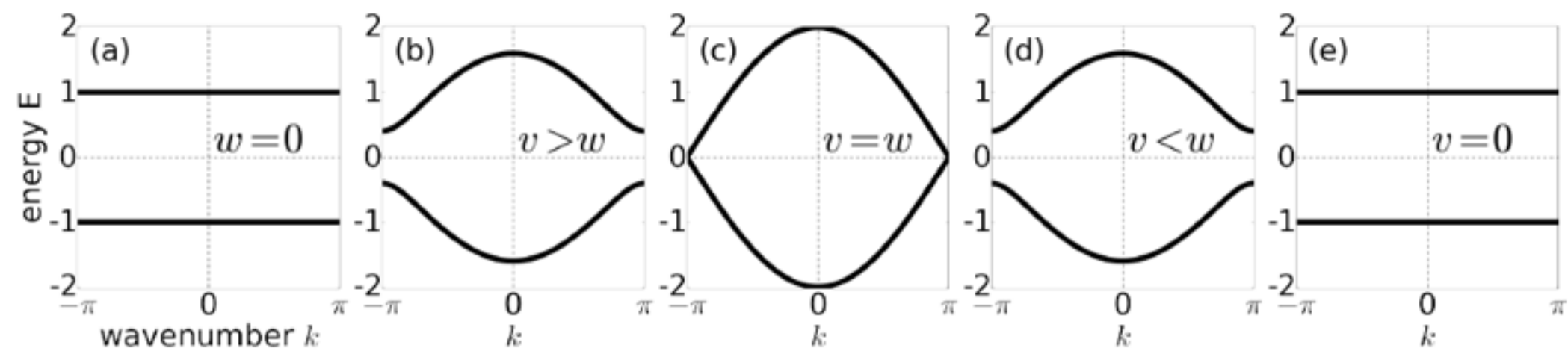
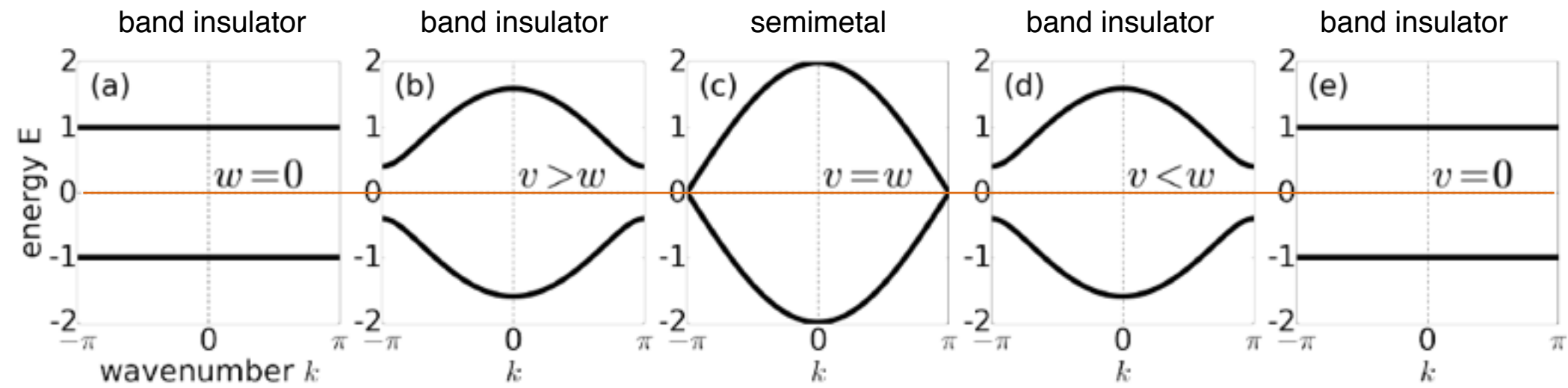
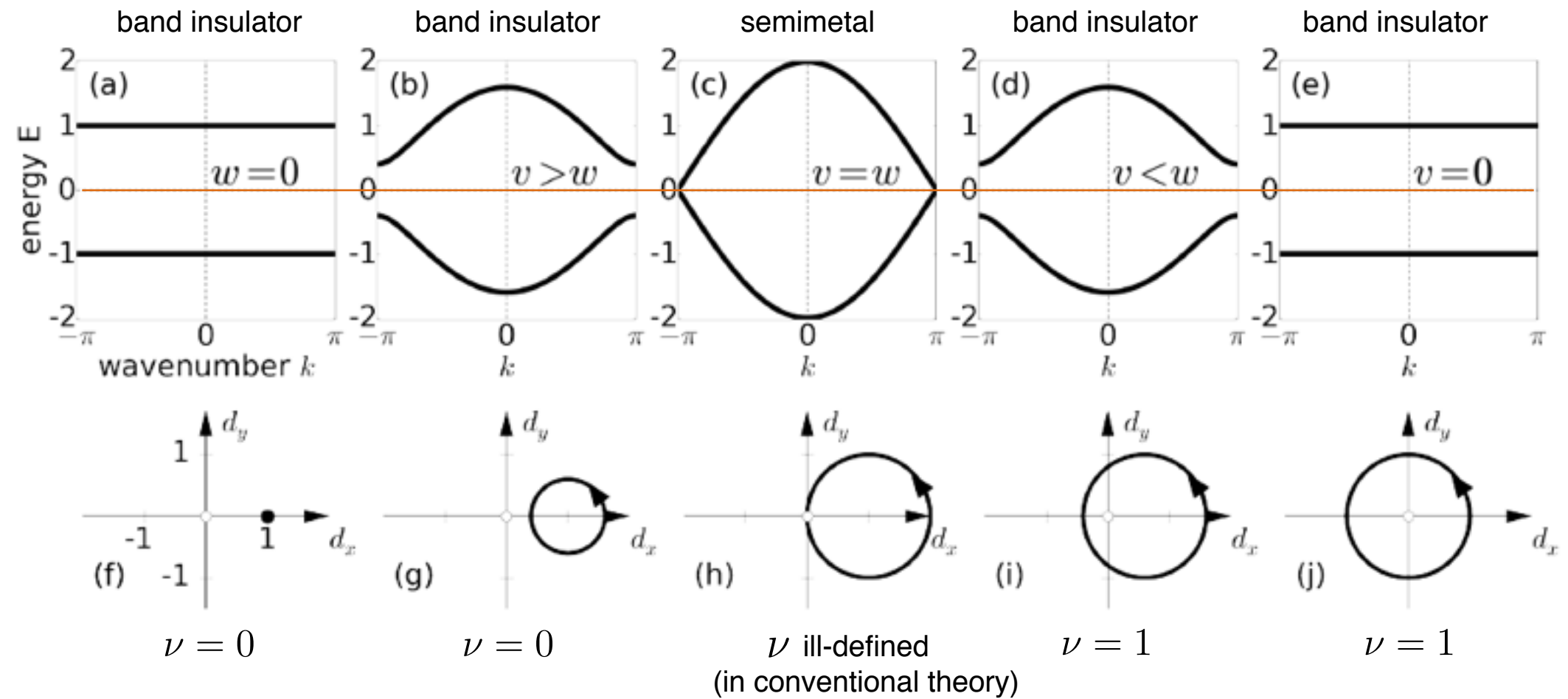




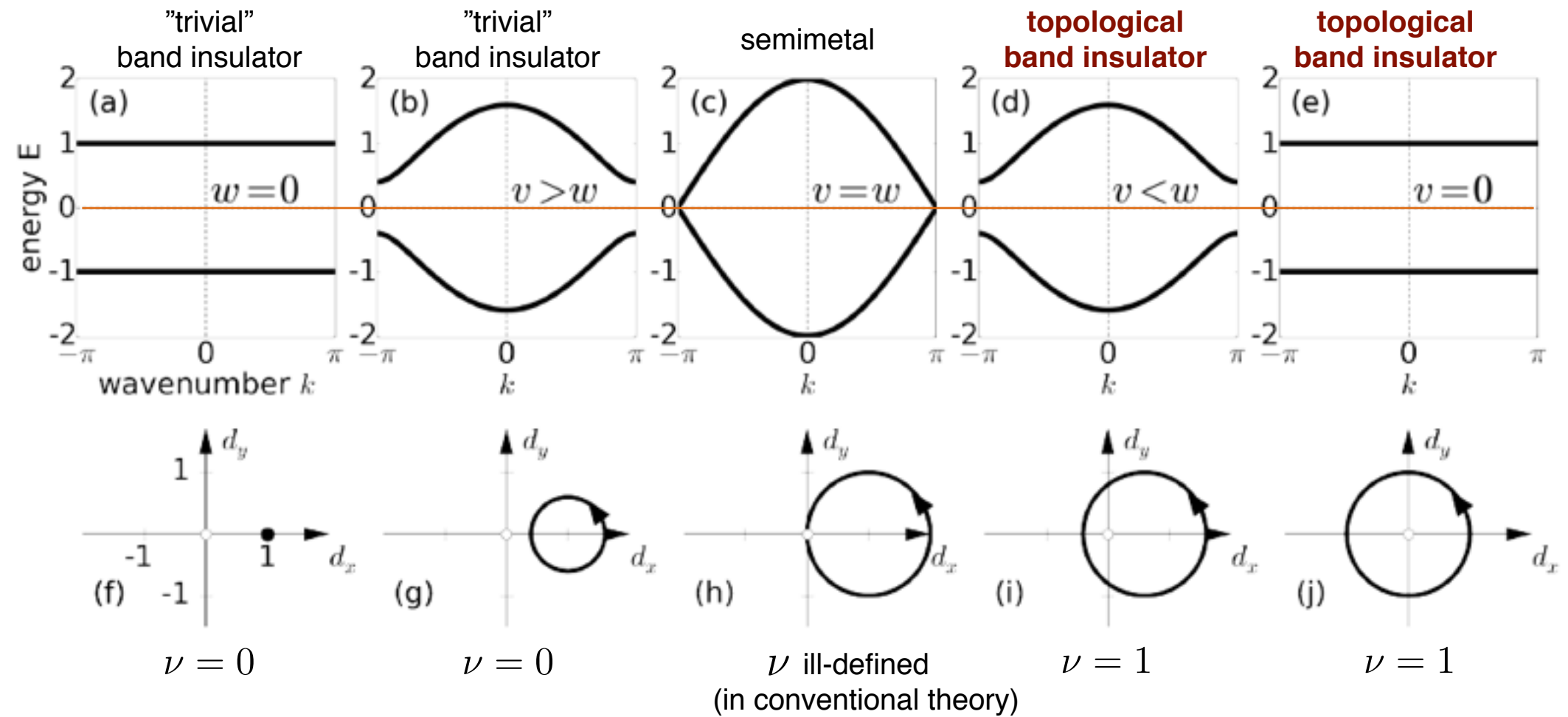
"half-filling": $E_{\text{Fermi}} = 0$



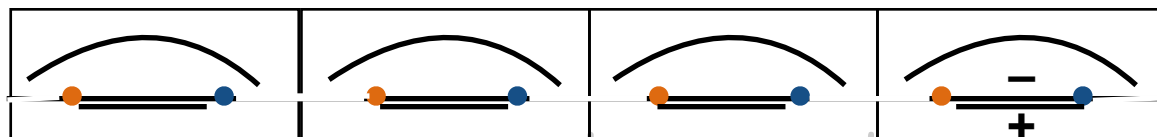
"half-filling": $E_{\text{Fermi}} = 0$



"half-filling": $E_{\text{Fermi}} = 0$



$$v \neq 0 \quad w = 0$$

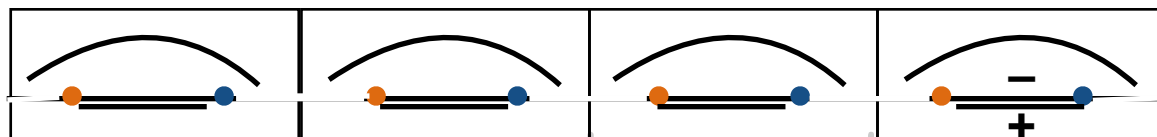


$$a = 1$$

$$\nu = 0$$

$$P = 0 \bmod e$$

$$v \neq 0 \quad w = 0$$



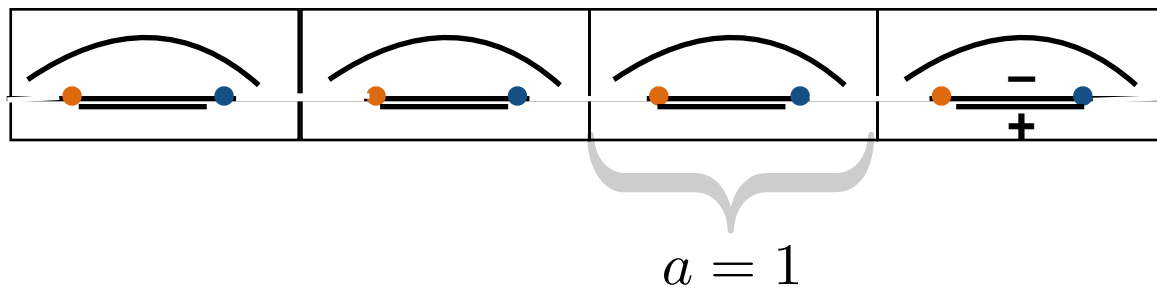
$$a = 1$$

$$\nu = 0$$

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"gauge fixing"

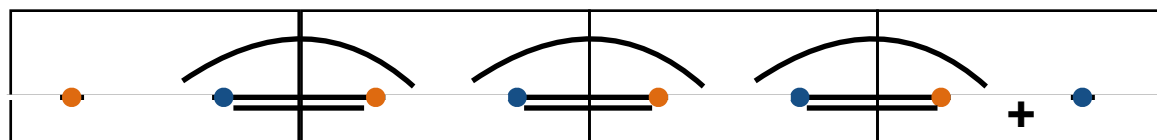
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$$\nu = 0$$

$$P = 0 \bmod e$$

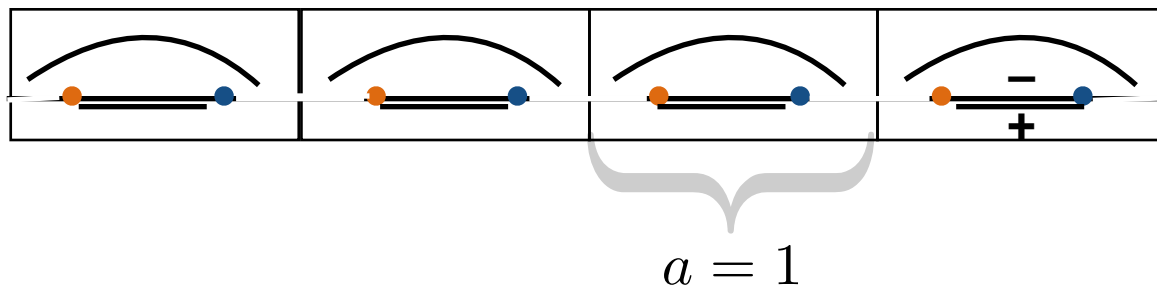
$$w \neq 0 \quad v = 0$$



$$\nu = 1$$

$$P = e/2 \bmod e$$

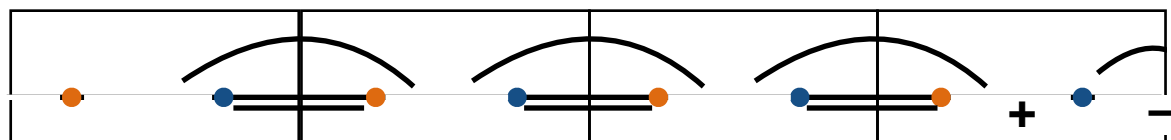
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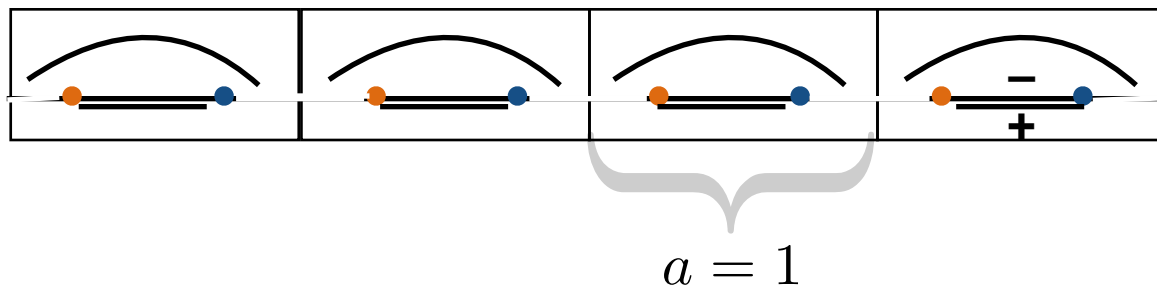
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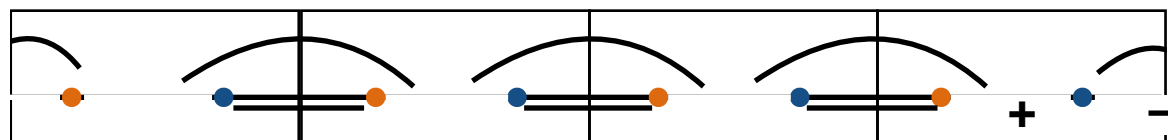
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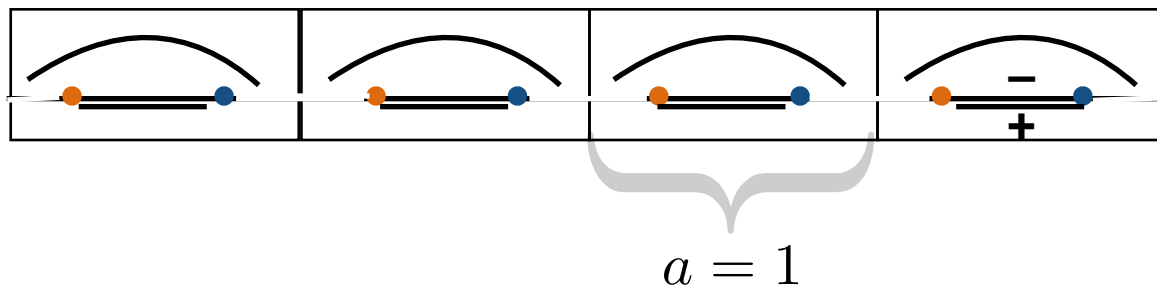
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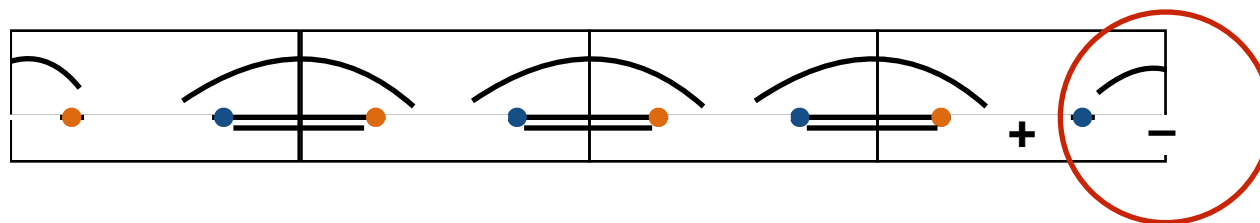
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$$P = 0 \bmod e$$

$$w \neq 0 \quad v = 0$$

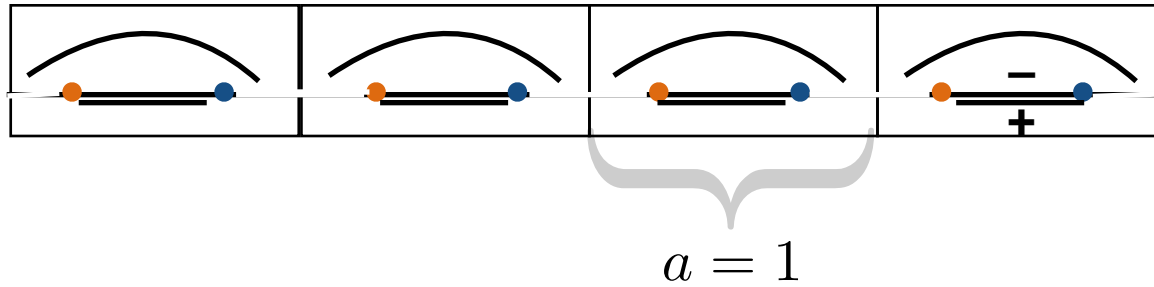


$$\nu = 1$$

$$P = e/2 \bmod e$$

localized zero-energy
boundary state

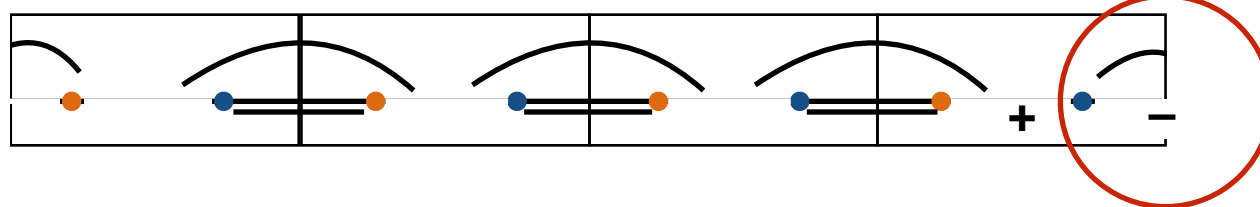
$$v \neq 0 \quad w = 0$$



$$\nu = 0$$

$$P = 0 \bmod e$$

$$w \neq 0 \quad v = 0$$

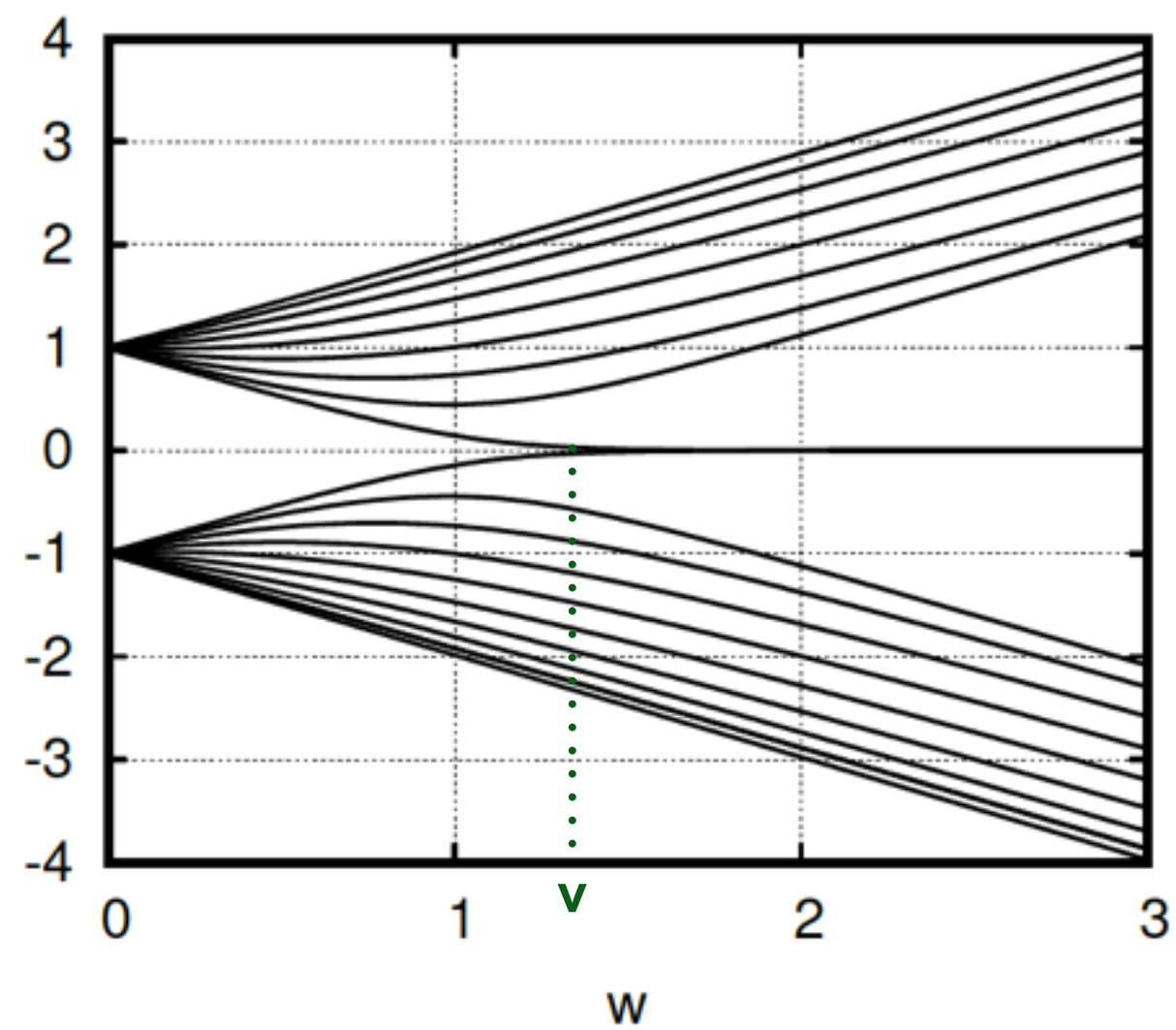


$$\nu = 1$$

$$P = e/2 \bmod e$$

localized zero-energy
boundary state

This boundary state is robust against perturbations which respect certain symmetries. This is THE hallmark of "symmetry-protected topological phases"!



”Periodic table” of symmetry-protected topological systems (noninteracting fermions)

class	T	C	S	0	1	2	3
A	0	0	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0
AIII	0	0	1	0	$\mathbb{Z}$	0	$\mathbb{Z}$
AI	+	0	0	$\mathbb{Z}$	0	0	0
BDI	+	+	1	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0
D	0	+	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0
DIII	−	+	1	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$
AII	−	0	0	$2\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$
CII	−	−	1	0	$2\mathbb{Z}$	0	$\mathbb{Z}_2$
C	0	−	0	0	0	$2\mathbb{Z}$	0
CI	+	−	1	0	0	0	$2\mathbb{Z}$

Anderson Localization (Schnyder, Ryu, Furusaki, Ludwig, 2008; 2009; 2010)

Topology (K-Theory) (Kitaev, 2009)

review: C.-K. Chiu *et al.*, Rev. Mod. Phys. **88**, 035005 (2016)

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example:
integer quantum Hall effect
topological invariant: *Chern number*

example:
SSH model
topological invariant: *winding number*

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example:
SSH model
topological invariant: *winding number*
(when allowing only perturbations that have **all** the symmetries of the unperturbed Hamiltonian!)

Anderson Localization (Schnyder, Ryu, Furusaki, Ludwig, 2008; 2009; 2010)

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D	0	+	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0
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C	0	−	0	0	0	$2\mathbb{Z}$	0
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another example:
SSH model with a perturbation that
only respects time-reversal symmetry
no topological phase!

Anderson Localization (Schnyder, Ryu, Furusaki, Ludwig, 2008; 2009; 2010)

Topology (K-Theory) (Kitaev, 2009)

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CII	−	−	1	0	$2\mathbb{Z}$	0	$\mathbb{Z}_2$
C	0	−	0	0	0	$2\mathbb{Z}$	0
CI	+	−	1	0	0	0	$2\mathbb{Z}$

By **adding** space group symmetries to T, C, and S (eg. mirror symmetry), a topologically trivial phase (like the D=1 All phase) may split into a trivial and a topologically nontrivial phase.
Crystalline topological insulators!

Anderson Localization (Schnyder, Ryu, Furusaki, Ludwig, 2008; 2009; 2010)

Topology (K-Theory) (Kitaev, 2009)

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